Derivation of the mean radial displacement, \bar{l}_{ρ}^{2D} , in 2 dimensions

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May 26, 2009

1 Introduction

In the following we derive an expression for the mean radial displacement, \bar{l}_{ρ}^{2D} , for diffusing 2D molecules in analogy to the well known expression for \bar{l}_r in three dimensions.

2 Derivation

Starting point is the solution of the diffusion equation in 2 dimensions

$$\frac{\partial c(\rho, t)}{\partial t} = D\nabla^2 c(\rho, t) \tag{1}$$

which, for a point source of M molecules released at the origin at time t = 0 can be shown to be

$$c(\rho,t) = \frac{M}{4\pi Dt} e^{-\frac{\rho^2}{4Dt}}$$
(2)

(see also Eqs. 3.1, 3.2 in [1]). Therefore, the probability density for a molecule being in a radial shell of thickness $d\rho$ around the origin is given by

$$p^{2D}(\rho,t) = \frac{(2\pi\rho)d\rho}{4\pi Dt} e^{-\frac{\rho^2}{4Dt}}$$
(3)

or, in normalized coordinates $\tilde{\rho} = \frac{\rho}{\lambda}$ with the normalization constant $\lambda = \sqrt{4Dt}$

$$p^{2\mathrm{D}}(\tilde{\rho},t) = 2e^{-\tilde{\rho}^2}(\tilde{\rho}d\tilde{\rho}) \quad . \tag{4}$$

Hence, the mean radial displacement of a 2D molecule is given by

$$\bar{l}_{\rho}^{\rm 2D} = 2\lambda \int_0^\infty \tilde{\rho}^2 e^{-\bar{\rho}^2} d\bar{\rho} \tag{5}$$

Since

$$\int_{0}^{\infty} t^{2n} e^{-at^{2}} dt = \frac{\Gamma(n+\frac{1}{2})}{2a^{n+\frac{1}{2}}}$$

$$\Rightarrow \int_{0}^{\infty} \tilde{\rho}^{2} e^{-\tilde{\rho}^{2}} d\tilde{\rho} = \frac{\Gamma(\frac{3}{2})}{2} = \frac{\sqrt{\pi}}{4}$$
(6)

and from this

$$\left| \bar{l}_{\rho}^{\rm 2D} = 2\lambda \frac{\sqrt{\pi}}{4} = \sqrt{\pi Dt} \right| \tag{7}$$

References

 R. Kerr, T.M. Bartol, B. Kaminsky, M. Dittrich, J.C.J. Chang, S. Baden, T.J. Sejnowski, and J.R. Stiles, Fast Monte Carlo Simulation Methods for Biological Reaction-Diffusion Systems in Solution and on Surfaces (2008), SIAM J. Sci. Comput., 30:3126-3149.